

Aletha de Witt AVN/Newton-fund 2017 Observational & Technical Training HartRAO





The HartRAO 26m telescope => equatorially mounted Cassegrain radio telescope

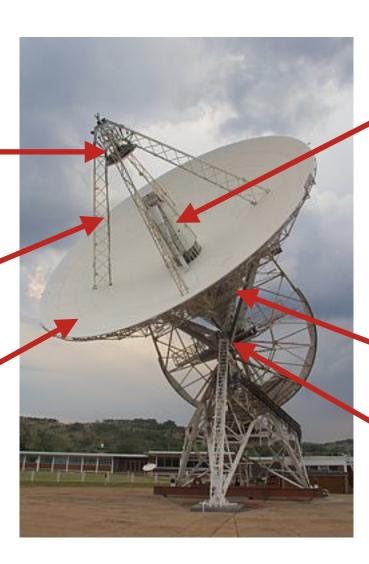
The **antenna reflectors** concentrate incoming E-M radiation into the focal point of the antenna

#### Secondary reflector

Sub-reflector (small reflector of hyperbolic curvature in front of the focus of the main reflector).

Sub-reflector support legs

**Primary reflector** 



Feed housing (feed horns receivers and support structure)

Converts E-M radiation in free space to electrical currents in a conductor.

#### 26 m telescope receivers (7):

1.6, 2.3, 5, 6.7, 8.4, 12.2 GHz 5 & 8.4 GHz **dual beam** new 22 GHz cooled receiver 15 GHz **dual beam** coming

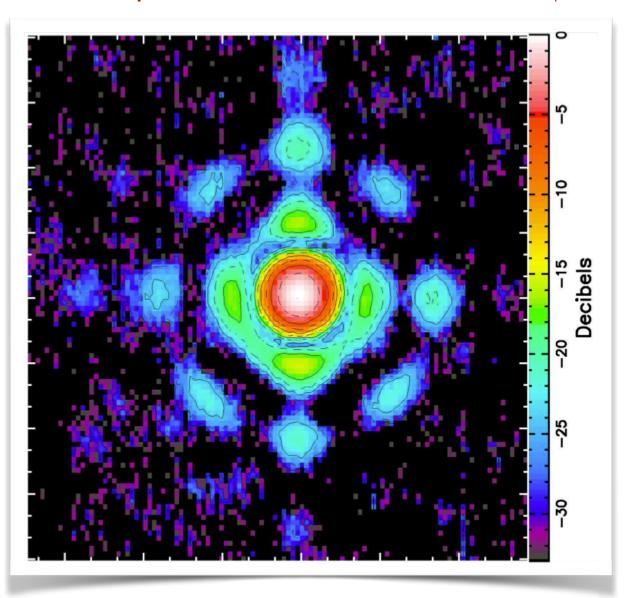
Deck Room

Local oscillator and mixers

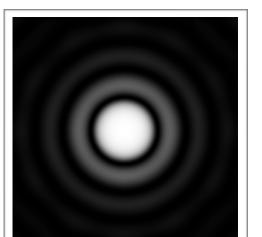
The antenna positioner
points the antenna at the desired location in the sky.



Actual beam pattern at 2300 MHz of the HartRAO 26m telescope

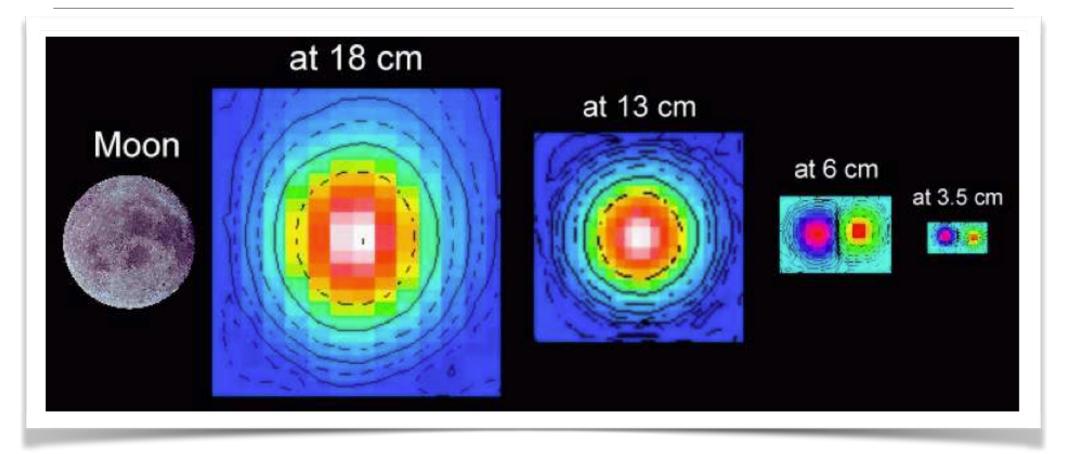


Diffraction pattern of a circular lens or reflector







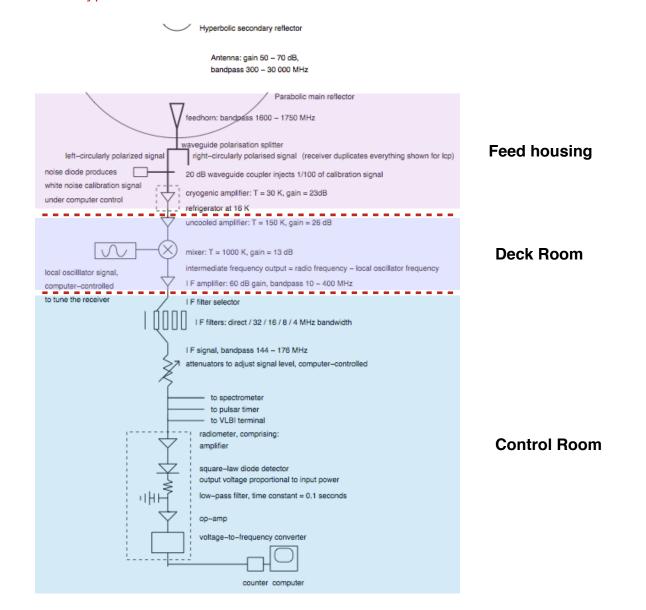


The **size of the main beam** of the 26-m telescope:

Depends on the **operating frequency/wavelength**. Here the actual observed main beam at four wavelengths are shown with the angular size of the Moon for comparison. Dual feeds on the 6 and 3.5 cm receivers produce two beams.



#### Signal chain: Main components of a typical microwave receiver and radiometer

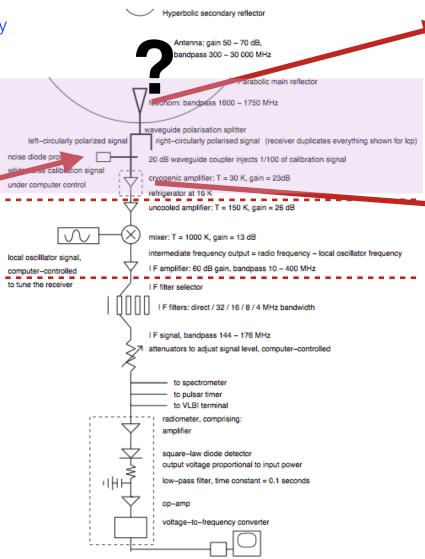




Signal chain: Main components of a typical microwave receiver and radiometer

**Incoming signal:** are very faint and noise like.

To calibrate the system a high stability **noise diode** injects a known noise signal which is split equally by a power divider between the LCP and RCP receiver chains.



counter computer

Feed horn and waveguide (to connect feed horn to first amplifier).
All incoming signals are split into
LCP & RCP by a hybrid
waveguide polarisation splitter feeding LCP to one receiver chain and RCP to the other.

Amplification to a
detectable level through a
low-noise amplifier.
Because the internal
noise in the amplifiers is
generally much larger than
the signal, specially designed
amplifiers that are
cryogenically cooled are
used to maximize sensitivity.

## Detecting Radio Emission from Space



Signal chain: Main components of a typical microwave receiver and radiometer

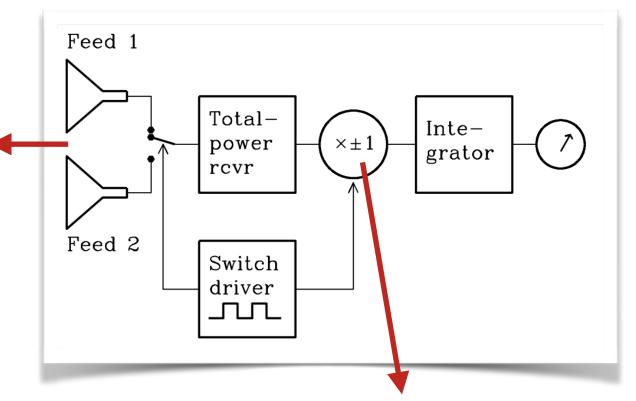
If **feed 1** is pointing at the source (angular size of source smaller than separation of the beams from the two feeds) then **feed 2** will point off-source but measure nearly the same sample of atmosphere in the near field.

#### **Dicke-switching:**

switching rapidly between two

#### identical feed horns

that are installed **East-West** next to each other on the telescope.



Output of receiver is **multiplied by +1** when receiver is connected to **feed 1 and by -1** when connected to **feed 2.** Fluctuations in atmospheric emission and drifts due to changes in receiver gain are canceled for frequencies below the switching rate.



Signal chain: Main components of a typical microwave receiver and radiometer

**Local oscillator** signal: computer–controlled to tune the receiver

To get the final output the IF signal is amplified, this time using an **IF Amplifier** 

Hyperbolic secondary reflector Antenna: gain 50 - 70 dB, bandpass 300 - 30 000 MHz Parabolic main reflecto feedhorn: bandpass 1600 - 1750 MHz waveguide polarisation splitter left-circularly polarized signal right-circularly polarised signal (receiver duplicates everything shown noise diode produces 20 dB waveguide coupler injects 1/100 of calibration signal white noise calibration signal cryogenic amplifier: T = 30 K, gain = 23dB under computer control refrigerator at 16 K uncooled amplifier: T = 150 K, gain = 26 dB mixer: T = 1000 K, gain = 13 dB intermediate frequency output = radio frequency - local oscillator frequency I F amplifier: 60 dB gain, bandpass 10 - 400 MHz computer-controlled F filters: direct / 32 / 16 / 8 / 4 MHz bandwidth I F signal, bandpass 144 - 176 MHz attenuators to adjust signal level, computer-controlled to pulsar time to VLBI terminal radiometer, comprising: amplifier square-law diode detector output voltage proportional to input power low-pass filter, time constant = 0.1 seconds op-amp voltage-to-frequency converter

counter computer

RF signal is **down converted** to a lower frequency in order to minimise signal losses in coaxial cable).

The **mixer** multiplies the RF signal with the **local oscillator signal.** The output signal that is used is the difference frequency component (RF - LO) of the product and is called the **intermediate frequency (IF).** 



Signal chain: Main components of a typical microwave receiver and radiometer

IF signal can be used unfiltered, or passed through 4, 8, 16 or 32-MHz bandwidth filters to exclude interference from external signals at some observing frequencies.

Voltage to frequency converter converts the signal to a square wave train (amplitude remains constant but the frequency is proportional to the DC voltage input).

These oscillations are then measured with a **counter** such that the count rate (in units of Hertz) is proportional to the original IF signal's power.

lyperbolic secondary reflector Antenna: gain 50 - 70 dB, bandpass 300 - 30 000 MHz Parabolic main reflector feedhorn: bandpass 1600 - 1750 MHz left-circularly polarized signal right-circularly polarised signal (receiver duplicates everything shown for lcp) 20 dB waveguide coupler injects 1/100 of calibration signal white noise calibration signal cryogenic amplifier: T = 30 K, gain = 23dB under computer control refrigerator at 16 K uncooled amplifier: T = 150 K, gain = 26 dB mixer: T = 1000 K, gain = 13 dB ntermediate frequency output = radio frequency I F amplifier: 60 dB gain, bandpass 10 - 400 MHz I F signal, bandpass 144 level, computer-controlled amplifier voltage-to-frequency converte counter computer

The **radiometer** is the basic instrument for measuring the power of the incoming signal. The simplest form of radiometer is the **"total power"** type shown

The signal is then detected by a **Square law detector** which converts the IF signal into an output DC voltage proportional to the input power.

Signals are loaded onto the Hart26m server in **FITS** (**Flexible Image Transport System**) format

# Detecting Radio Emission from Space



- The antenna needs to be calibrated to convert the signal amplitude in units of Hertz to units of Antenna Temperature in Kelvins [K], as it is the standard physically meaningful scale used with most radio analysis techniques.
- The output signal from the radiometer is proportional to the **T**<sub>sys</sub>, from which we can extract the **T**<sub>A</sub>.

$$T_{sys} = T_{Bcmb} + T_A + T_{at} + T_{wv} + T_g + T_R \text{ [K]}$$

- Prior to each drift scan, the noise diode injects a noise signal with a known temperature and this is used to calibrate the antenna.
- Comparing the noise diode's temperature to its count rate can derive a conversion factor [K/Hz] to convert from counts (Hz) to antenna temp (K).



The aperture efficiency can be obtained at each frequency;

$$\epsilon_{ap} = rac{A_e}{A_p} \;\;\;$$
 => max achievable aperture efficiency ~ 0.64

- Ae is the effective aperture (collecting area) and Ap is the physical collecting area, obtainable from the known diameter of the telescope (25.9 m for the HartRAO 26m telescope).
- The source flux density S, is the product of the brightness and source solid angle

$$h\nu << kT, \ B = \frac{2kT}{\lambda^2} \ [\text{W m}^{-2} \ \text{Hz}^{-1} \ \text{sr}^{-1}]$$
 
$$S = \frac{2kT\Omega_s}{\lambda^2} \ [\text{W m}^{-2} \ \text{Hz}^{-2}]$$

Remember !!! 1 Jy = 
$$10^{-26}$$
 [W m<sup>-2</sup> Hz<sup>-2</sup>]

## Radiometer Equation



- Radio Astronomers like to think of their telescopes as resistors .....
  - .. and when you put power into a resistor
  - ... it heats up

$$h\nu << kT, \ B = \frac{2kT}{\lambda^2} \ [Wm^{-2}Hz^{-2}Sr^{-2}]$$

Rayleigh-Jeans Law holds all the way through the radio regime for any reasonable temperature.

The question is: what flux density is received by your antenna



$$\int Bd\Omega = S \left[ Wm^{-2}Hz^{-2} \right]$$

Remember !!! 1 Jy = 
$$10^{-26}$$
 [W m<sup>-2</sup> Hz<sup>-2</sup>]

## Radiometer Equation



 Now lets look at the power that we actually received by the antenna at a given frequency .....

... we integrate the flux density over the area of the antenna

$$\int SdA = P \left[ WHz^{-2} \right]$$

- Now the antenna theorem states:  $A_e\Omega=\lambda^2$
- Lets go one step back from power (without using fancy integration)
  ... what we effectively just did was ...

$$B\Omega A_e \approx \frac{2kT}{\lambda^2}\Omega A_e$$
  $SA_e = 2kT$ 



- We obtain the true flux density of the source by summing the antenna temperatures measured in RCP and LCP.
- The total intensity is the sum of what is received in each polarisation.

$$S = \frac{k(T_{Alcp} + T_{Arcp})K_s}{A_e} \times 10^{26} \text{ [Jy]}$$

- To obtain the true flux density S we introduce a size correction factor Ks. For sources that are very small compared to the beam size, Ks = 1, but the correction must be taken into account if the source size is a significant fraction of the beam size.
- We can only calculate the source flux density if we know the effective aperture (collecting area) at the frequency being used, so we rewrite above equation and substitute the constants, to give;

$$A_e = \frac{1380(T_{Alcp} + T_{Arcp})K_s}{S_o} \text{ [m}^2]$$



- It is important to note that the **flux density** of a radio source is **intrinsic** to it, and the same flux density should be measured by any properly calibrated telescope. However the antenna temperatures measured for the same emitter by different telescopes will be proportional to their effective collecting areas.
- We can now calibrate the telescope at each frequency of interest. We can carry out scans of **standard calibrator sources** (Ott et al. 1994) and measure the peak antenna temperature in each polarisation.



- For convenience, we often refer to the **Point Source Sensitivity** (*PSS*), which is the number of Kelvins of antenna temperature per polarisation, obtained per Jansky of source flux density. This is also known as the '**DPFU**' or '**Degrees per Flux Unit**'.
- For the HartRAO 26 m telescope the *PSS* is typically about 5 Jy/Kelvin per polarisation. The **PSS** in each polarisation is simple to determine experimentally from the measured T<sub>A</sub> of calibrator sources of known flux density. **NB: unpolarised sources => half the total flux density is received in each polarisation.**

$$PSS_{lcp} = \frac{(S/2)}{K_s T_{Alcp}}$$
 and  $PSS_{rcp} = \frac{(S/2)}{K_s T_{Arcp}}$  [Jy K<sup>-1</sup> per polarisation]

 Theoretically the values for the two polarisations should be the same; in practise there is always a small difference between them, and data from each polarisation should be corrected using the value appropriate for that polarisation.

## Detecting Radio Emission from Space



- Simplest way to measure the intensity of a **compact source** in the sky, i.e. one that has an angular size much smaller than the beam, is to use an observing method called a **drift scan**.
- The output of the radiometer will be the convolution of the antenna beam pattern with the brightness distribution of the source.
- If the source is compact, the output from the radiometer during the scan is effectively an **east-west cross-section of the beam** of the telescope.

## Detecting Radio Emission from Space



An example of a drift scan

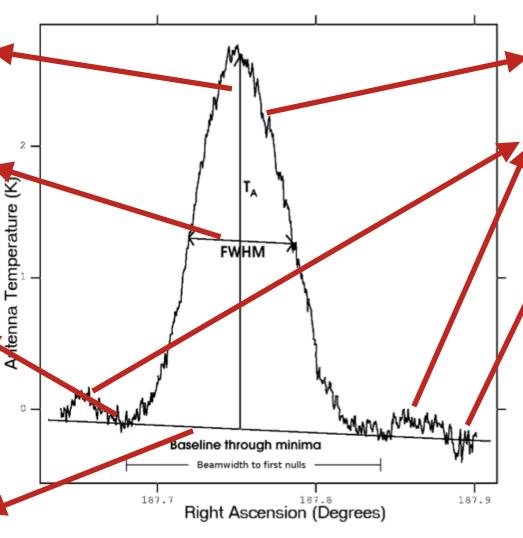
We measure the height above the baseline at the centre of the beam to get the **antenna temperature.** 

We can also measure the **FWHM** if it is an unresolved source.

Looking at the **minima** across the scan we can see a **slow drift** in the signal level.

This could be due to changing atmospheric conditions or a slow change in the gain of the receiver.

We need to establish the slope between the first nulls by drawing a line between them.



The passage of the **main beam** across the radio source is in the centre

The first **side lobes** are seen weakly on each side.

The **noise** is clearly visible.

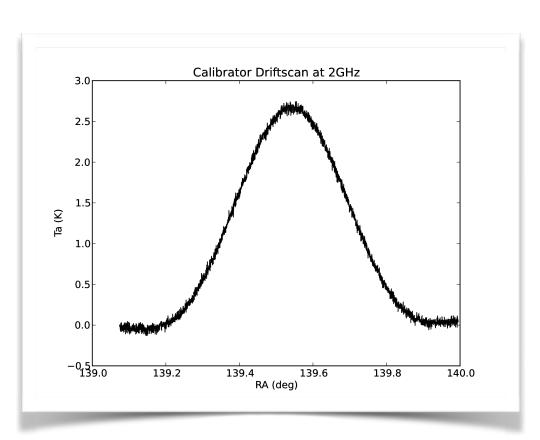
If the source is a calibrator, the *PSS* in this polarisation is obtained from the flux density S at the observing frequency (Ott et al. 1994).

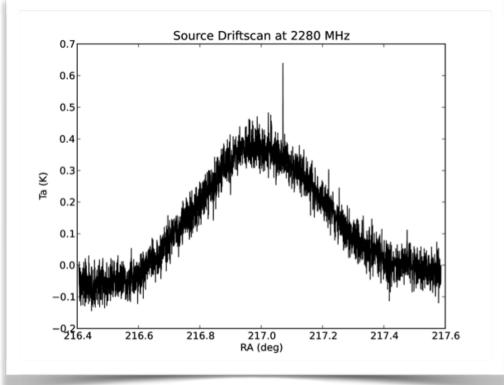
Then we can find the flux density of unknown sources from their antenna temperature.

## Monitoring of Active Galactic Nuclei



#### HartRAO 26 m telescope, drift scans => raw data



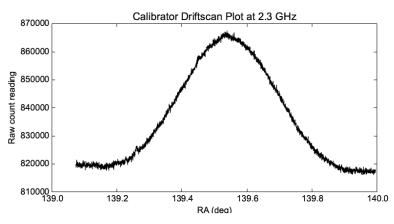


Calibrator: Hydra A Image Credit: Pfesesani Nemanashi, Mike Gaylard Source: J1427-4206 / PKS 1424-418Image Credit: Pfesesani Nemanashi, Mike Gaylard

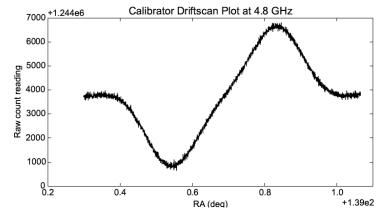
## Monitoring of Active Galactic Nuclei



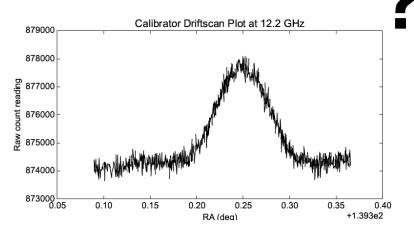
#### HartRAO 26 m telescope, drift scans => raw data



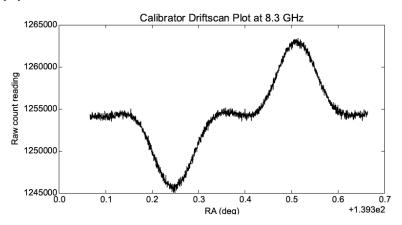
(a) Drift scan pattern output at 2.3 GHz.



(c) Drift scan pattern output at 4.8 GHz.



(b) Drift scan pattern output at 12.2 GHz.

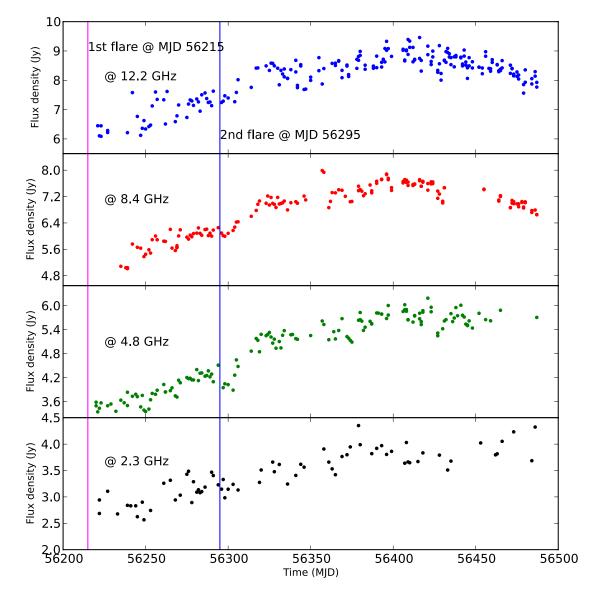


(d) Drift scan pattern output at 8.3 GHz.

# Monitoring of Active Galactic Nuclei



Monitoring of J1427-4206 - HartRAO 26 m



### Data Reduction



#### 3C218/Hydra A

J2000 coordinates RA:09h18m05.67s Dec: -12°05m44.0s

#### **Equation 1**

$$A_e = \frac{1380(T_{Alcp} + T_{Arcp})K_s}{S_o}$$
 [m<sup>2</sup>]

Equation 2 
$$\epsilon_{ap} = \frac{A_e}{A_p}$$

$$A_p = \frac{\pi}{4}D^2$$
  $D = 25.9 \text{ m}$ 

#### **Equation 3**

$$PSS_{lcp} = \frac{(S/2)}{K_s T_{Alcp}}$$
 and  $PSS_{rcp} = \frac{(S/2)}{K_s T_{Arcp}}$  [Jy K<sup>-1</sup> per polarisation]