Radio Continuum Observations

Aletha de Witt DARA-AVN May 2019 Observational & Technical Training HartRAO

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The HartRAO 26m telescope => equatorially mounted Cassegrain radio telescope

The **antenna reflectors** concentrate incoming E-M radiation into the focal point of the antenna

Secondary reflector

Sub-reflector (small reflector of hyperbolic curvature in front of the focus of the main reflector).

> Sub-reflector support legs

Primary reflector



Feed housing (feed horns receivers and support structure)

Converts E-M radiation in free space to electrical currents in a conductor. **26 m telescope receivers (7):** 1.6, 2.3, 5, 6.7, 8.4, 12.2 GHz 5 & 8.4 GHz **dual beam**

> Deck Room Local oscillator and mixers

Antenna positioner The antenna positioner points the antenna at the desired location in the sky.



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Actual beam pattern at 2300 MHz of the HartRAO 26m telescope



Diffraction pattern of a circular lens or reflector



The beam cross-section of an antenna



Angle from Beam Axis [wavelength/diameter] (radians)





The size of the main beam of the 26-m telescope: depends on the operating frequency/wavelength.

Here the actual observed main beam at four wavelengths are shown with the angular size of the Moon for comparison. Dual feeds on the 6 and 3.5 cm receivers produce two beams.



Dual feeds on the 6 and 3.5 cm receivers produce two beams. Identical feed horns that are installed East-West next to each other on the telescope.



Output of receiver is multiplied by +1 when receiver is connected to feed 1 and by
-1 when connected to feed 2. Fluctuations in atmospheric emission and drifts due to changes in receiver gain are canceled for frequencies below the switching rate.





























Factors reducing the **aperture efficiency** (0.80, 0.75, 0.64)





Pointing accuracy



RFI





Atmosphere





- The antenna needs to be calibrated to convert the signal amplitude in units of Hertz to units of Antenna Temperature in Kelvins [K], as it is the standard physically meaningful scale used with most radio analysis techniques.
- The output signal from the radiometer is proportional to the **T**_{sys}, from which we can extract the **T**_A.

$$T_{sys} = T_{Bcmb} + T_A + T_{at} + T_{wv} + T_g + T_R [K]$$

- Prior to each drift scan, the noise diode injects a noise signal with a known temperature and this is used to calibrate the antenna.
- Comparing the noise diode's temperature to its count rate can derive a conversion factor [K/Hz] to convert from counts (Hz) to antenna temp (K).



- Unlike line emission, we do not have a specific frequency to tune our radio receiver to when we are doing continuum observations. Ideally, to gain the optimal amount of information from our source, we'd like to observe over as wide a range of frequencies as possible. Because of the design of radio correlators and receivers, we are usually limited to some set frequency range.
- So, for example, you might be interested in the total flux from all radiation in your source between 2.5 cm and 3.5 cm, which you then call the "3cm continuum flux density". Thus you'll often see terms like the "7 mm continuum" and the "6 cm continuum".
- To get a full spectrum of your source across the entire radio range would require you to observe across a range of radio bands. This must be done either (i) at different times, (ii) at different telescopes, or (iii) at telescopes with dual or multi-frequency observing capacity. In reality this is very rare.



• The aperture efficiency can be obtained at each frequency;

$$\epsilon_{ap} = \frac{A_e}{A_p}$$
 => max achievable aperture efficiency ~ 0.64

- Ae is the effective aperture (collecting area) and Ap is the physical collecting area, obtainable from the known diameter of the telescope (25.9 m for the HartRAO 26m telescope).
- The source flux density S, is the product of the brightness and source solid angle

$$h\nu \ll kT, \ B = \frac{2kT}{\lambda^2} \ [W m^{-2} Hz^{-1} sr^{-1}]$$

 $S = \frac{2kT\Omega_s}{\lambda^2} \ [W m^{-2} Hz^{-2}]$

Remember !!! 1 Jy = 10^{-26} [W m⁻² Hz⁻²]



- Now the antenna theorem states: $A_e\Omega=\lambda^2$
- Lets go one step back from power (without using fancy integration) ... what we effectively just did was ...



 $SA_e = 2kT$



- We obtain the true flux density of the source by summing the antenna temperatures measured in RCP and LCP.
- The total intensity is the sum of what is received in each polarisation.

$$S = \frac{k(T_{Alcp} + T_{Arcp})K_s}{A_e} \times 10^{26} \text{ [Jy]}$$

- To obtain the true flux density S we introduce a size correction factor Ks. For sources that are very small compared to the beam size, *K*s = 1, but the correction must be taken into account if the source size is a significant fraction of the beam size.
- We can only calculate the source flux density if we know the effective aperture (collecting area) at the frequency being used, so we rewrite above equation and substitute the constants, to give;

$$A_e = \frac{1380(T_{Alcp} + T_{Arcp})K_s}{S_o} \, [\text{m}^2]$$



- It is important to note that the **flux density** of a radio source is **intrinsic** to it, and the same flux density should be measured by any properly calibrated telescope. However the antenna temperatures measured for the same emitter by different telescopes will be proportional to their effective collecting areas.
- We can now calibrate the telescope at each frequency of interest. We can carry out scans of **standard calibrator sources** (Ott et al. 1994) and measure the peak antenna temperature in each polarisation.



- For convenience, we often refer to the Point Source Sensitivity (PSS), which is the number of Kelvins of antenna temperature per polarisation, obtained per Jansky of source flux density. This is also known as the 'DPFU' or 'Degrees per Flux Unit'.
- For the HartRAO 26 m telescope the PSS is typically about 5 Jy/Kelvin per polarisation. The PSS in each polarisation is simple to determine experimentally from the measured TA of calibrator sources of known flux density. NB: unpolarised sources => half the total flux density is received in each polarisation.

$$PSS_{lcp} = \frac{(S/2)}{K_s T_{Alcp}}$$
 and $PSS_{rcp} = \frac{(S/2)}{K_s T_{Arcp}}$ [Jy K⁻¹ per polarisation]

• Theoretically the values for the two polarisations should be the same; in practise there is always a small difference between them, and data from each polarisation should be corrected using the value appropriate for that polarisation.



- For continuum emission:
- Simplest way to measure the intensity of a **compact source** in the sky, i.e. one that has an angular size much smaller than the beam, is to use an observing method called a **drift scan**.
- The output of the radiometer will be the convolution of the antenna beam pattern with the brightness distribution of the source.
- If the source is compact, the output from the radiometer during the scan is effectively an **east-west cross-section of the beam** of the telescope.
- This bandwidth is defined as B = f2 f1 where f1 is the lower operating frequency and f2 is the upper operating frequency

Thank You

Contact Details

Aletha de Witt alet@hartrao.ac.za